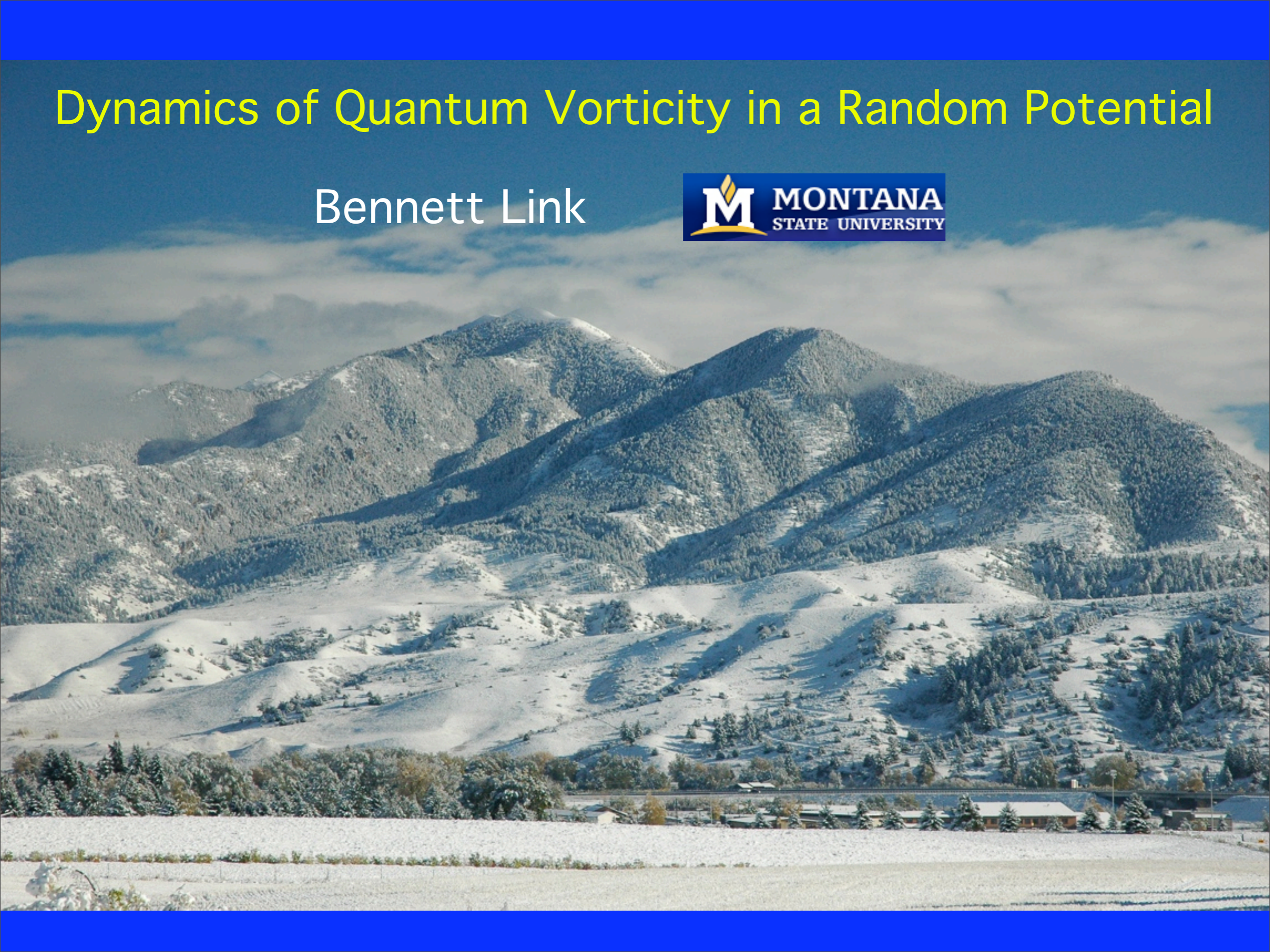


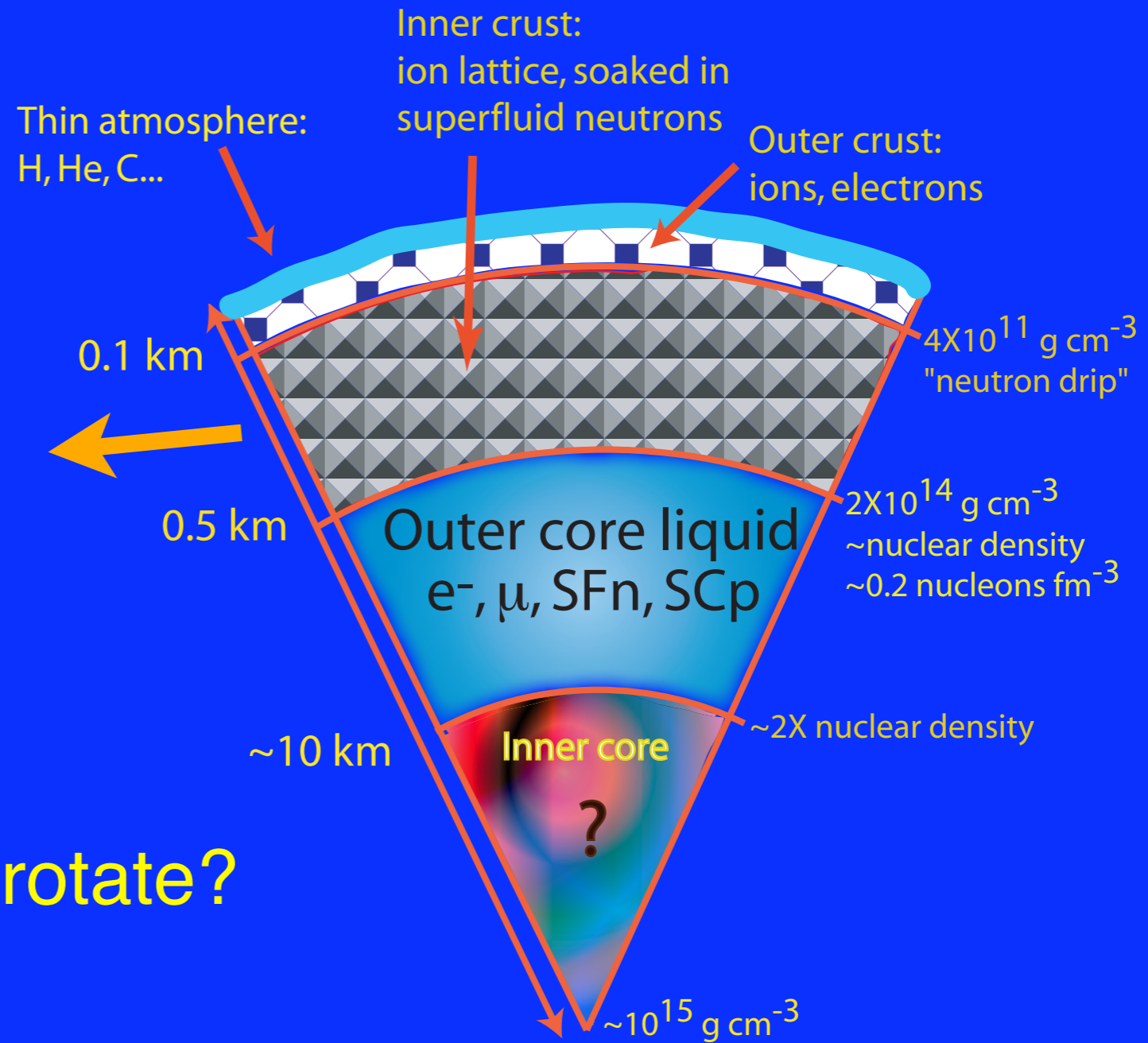
Dynamics of Quantum Vorticity in a Random Potential

Bennett Link



Neutron star consists of distinct superfluids

The inner-crust solid coexists with a neutron superfluid.



How does this SF rotate?

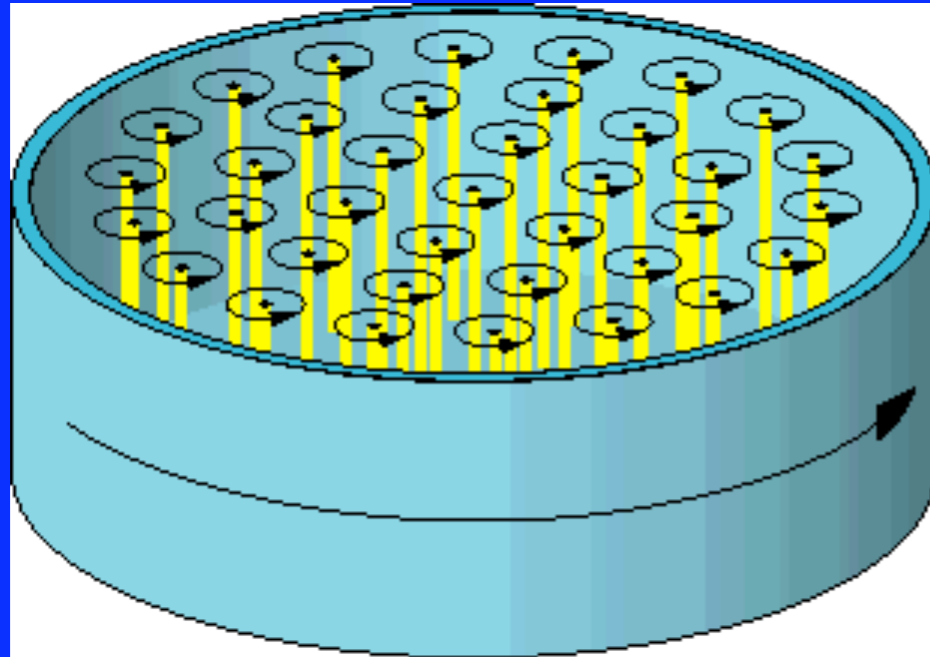
Why worry about SF rotational dynamics?

To develop a theory of NS “seismology”

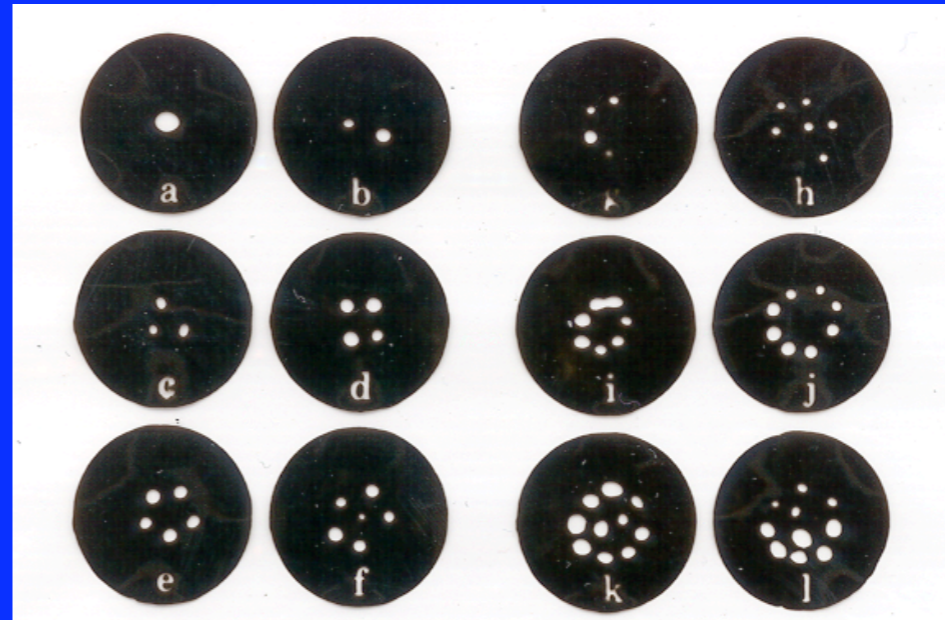
Some observed modes in NSs,
phenomena to explain/understand:

- Spin jumps (glitches).
- Precession (“wobble”, nutation).
- Stochastic spin variations (timing noise).
- Crust shear modes excited in magnetar flares.

The neutron SF's rotation



Rotating superfluid He



- Distribution of vortices determines SF angular momentum.
- Coupling between container (the crust) and the SF is determined by forces on the vortices.

changes in SF angular momentum
and SF torques on the crust



rearrangement of
vortex array



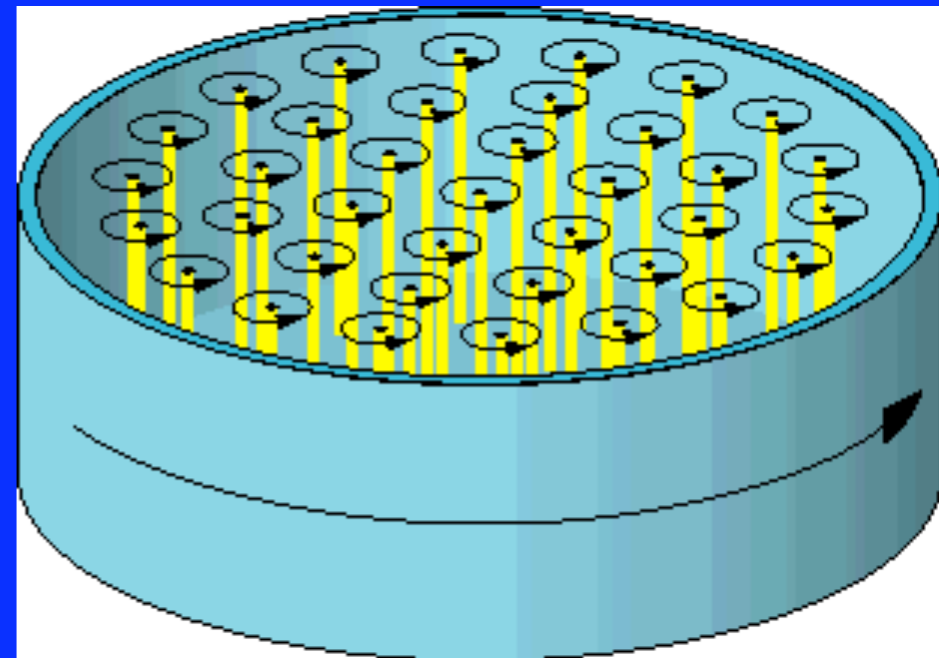
vortex mobility

How mobile are the inner-crust vortices in the lattice, and,
how is spin dynamics affected?

Dynamical considerations

Two velocities (in the frame of the container):

1. Velocity of the net SF flow (determined by overall vortex distribution)
2. Velocity of a vortex.



Vortex mobility (velocity) is determined by forces...

The forces

- A vortex interacts with nuclei, $E_{vn} \sim 1$ MeV.
(Donati & Pizzochero 06; Avogadro et al. 07)
- A vortex has self-energy (“tension”).
- There is a lift force (Magnus force) on a moving vortex segment.
- There is dissipation force associated with irreversible energy transfer to the lattice.

The forces

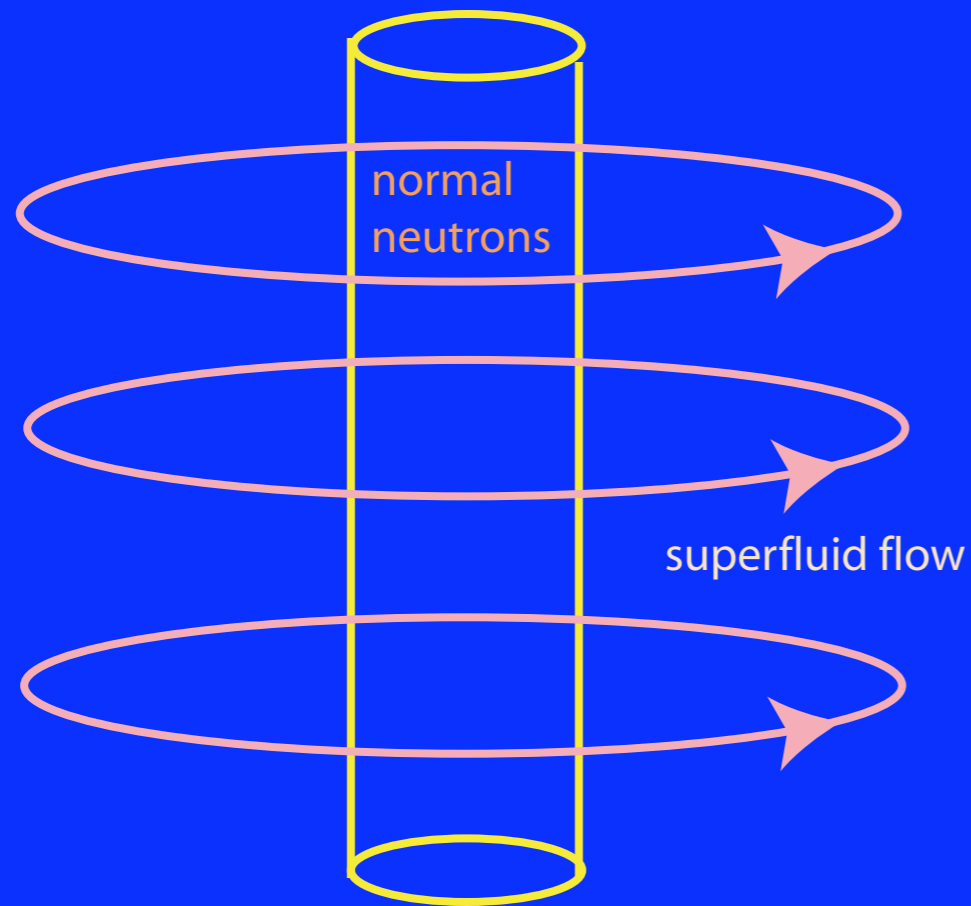
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Vortex/nucleus interaction due to
density dependence of the SF gap

The forces

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- A vortex has self-energy (“tension”).
- There is a lift force (Magnus force) on a moving vortex segment.
- There is dissipation force associated with irreversible energy transfer to the lattice.

A vortex has a large self-energy



$$\oint \mathbf{v}_s \cdot d\mathbf{l} = \kappa \equiv h/2m_n$$

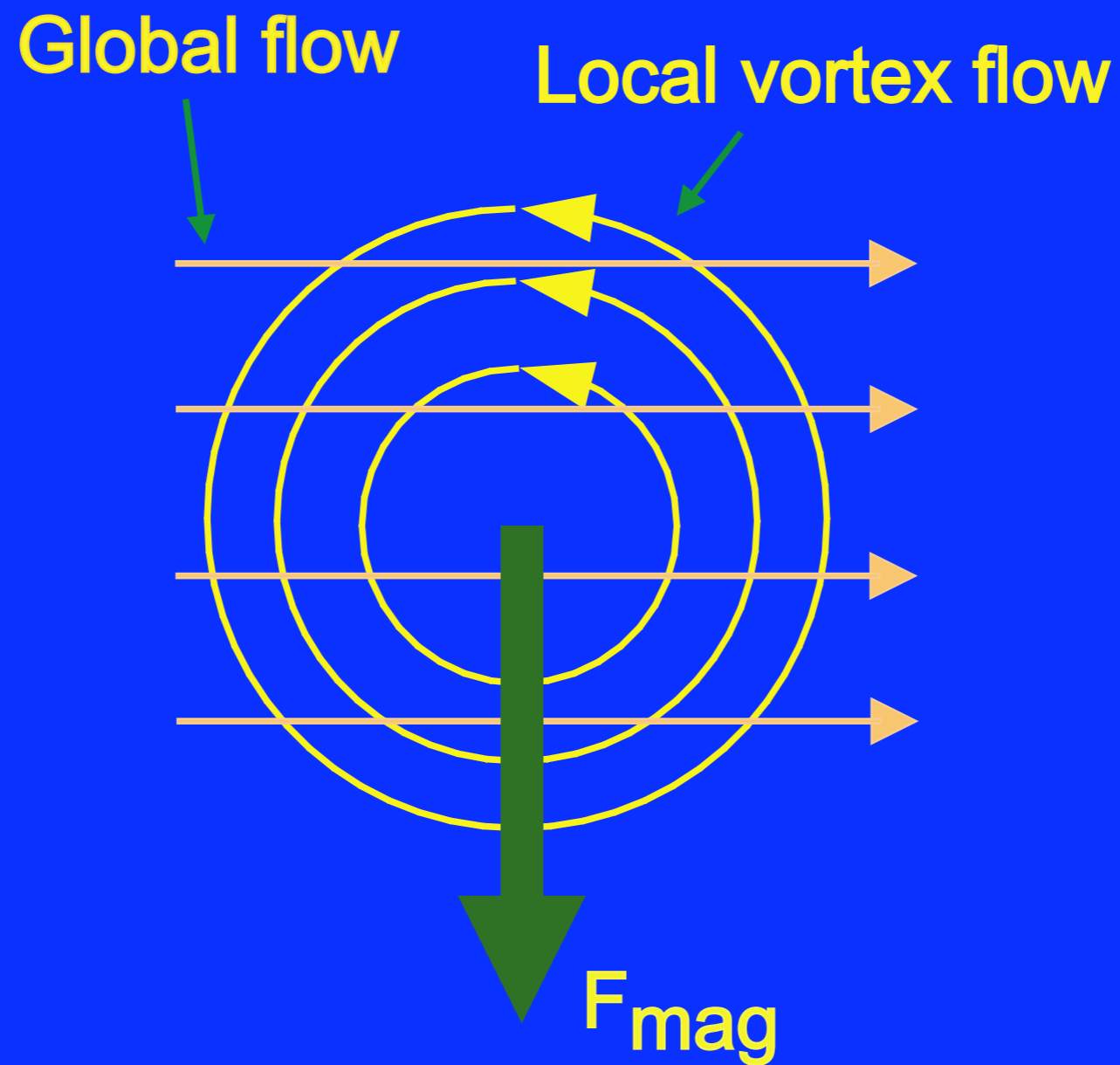
“Tension”

$$\frac{E_k}{L} \equiv T_v \simeq 10 \text{ MeV fm}^{-1}$$

The forces

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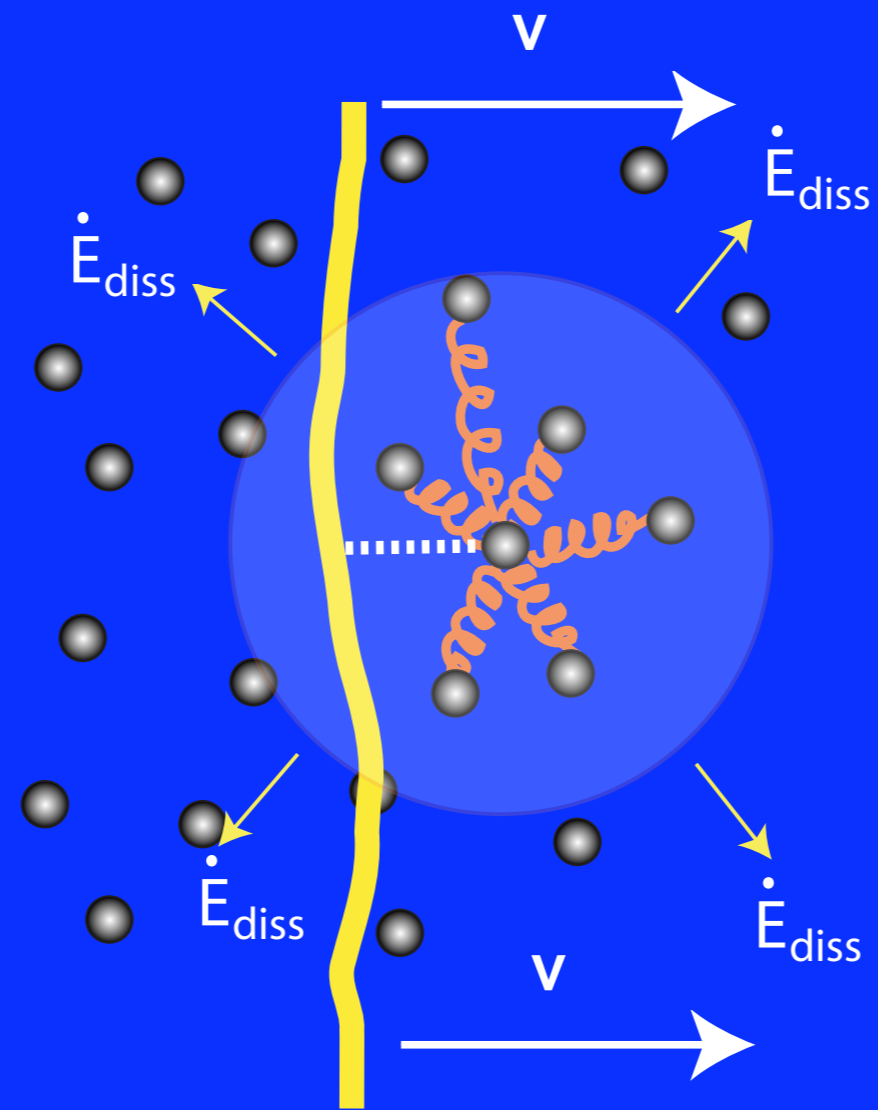
The Magnus force



The forces

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Vortex motion couples to phonons



Equations of motion

(Sonin 87; Link 08)

The diagram shows the equation of motion for a vortex, with several terms labeled by arrows:

- tension**: points to T_v
- vortex displacement**: points to $\frac{\partial^2 \mathbf{u}}{\partial z^2}$
- SF density**: points to ρ_s
- vorticity vector**: points to $\boldsymbol{\kappa}$
- ambient SF flow**: points to \mathbf{v}_s
- non-dissipative lattice force**: points to \mathbf{f}_0
- drag term**: points to $\eta \frac{\partial \mathbf{u}}{\partial t}$

$$T_v \frac{\partial^2 \mathbf{u}}{\partial z^2} + \rho_s \boldsymbol{\kappa} \times \left(\frac{\partial \mathbf{u}}{\partial t} - \mathbf{v}_s \right) + \mathbf{f}_0 - \eta \frac{\partial \mathbf{u}}{\partial t} = 0$$

Properties

$$T_v \frac{\partial^2 \mathbf{u}}{\partial z^2} + \rho_s \boldsymbol{\kappa} \times \left(\frac{\partial \mathbf{u}}{\partial t} - \mathbf{v}_s \right) + \mathbf{f}_0 - \eta \frac{\partial \mathbf{u}}{\partial t} = 0$$

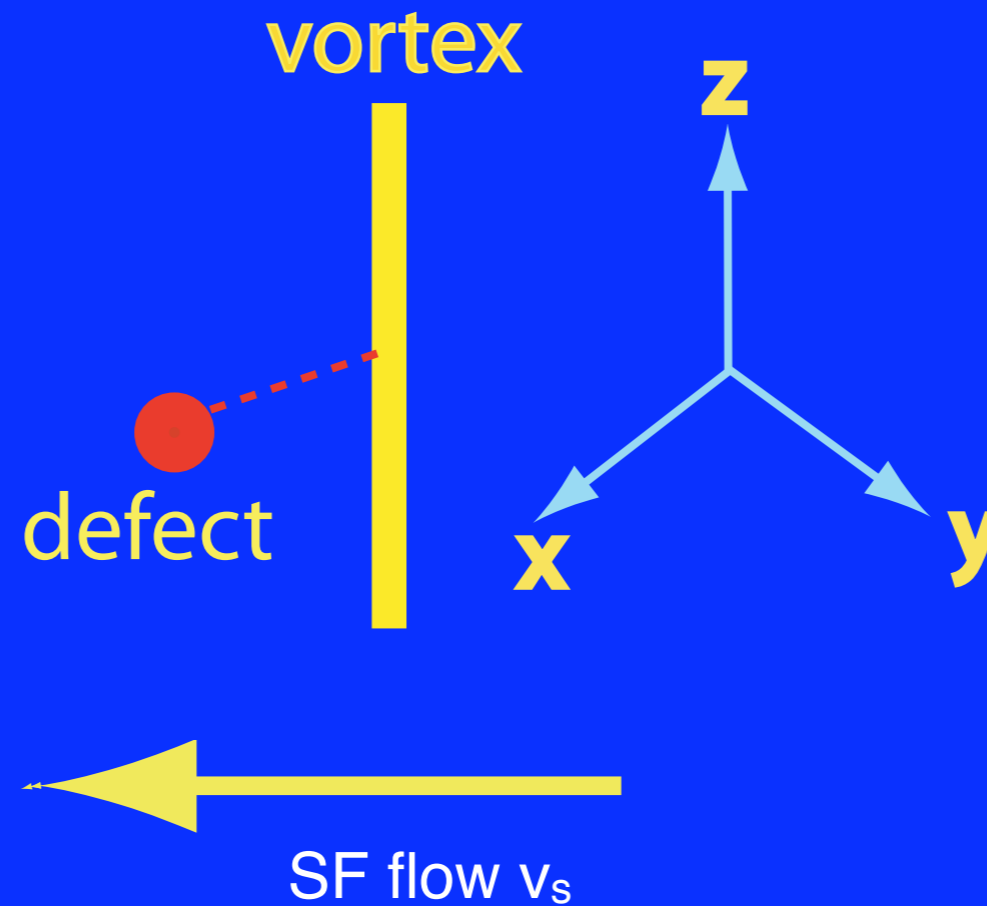
- Diffusive equation. No inertial term. Vortex supports helical waves, “kelvin modes”.
- Vortex moves transverse to an applied force.
- Velocity (not acceleration) is determined by applied forces. State of motion depends on position.

How does a vortex move
through the lattice?

Simpler problem:

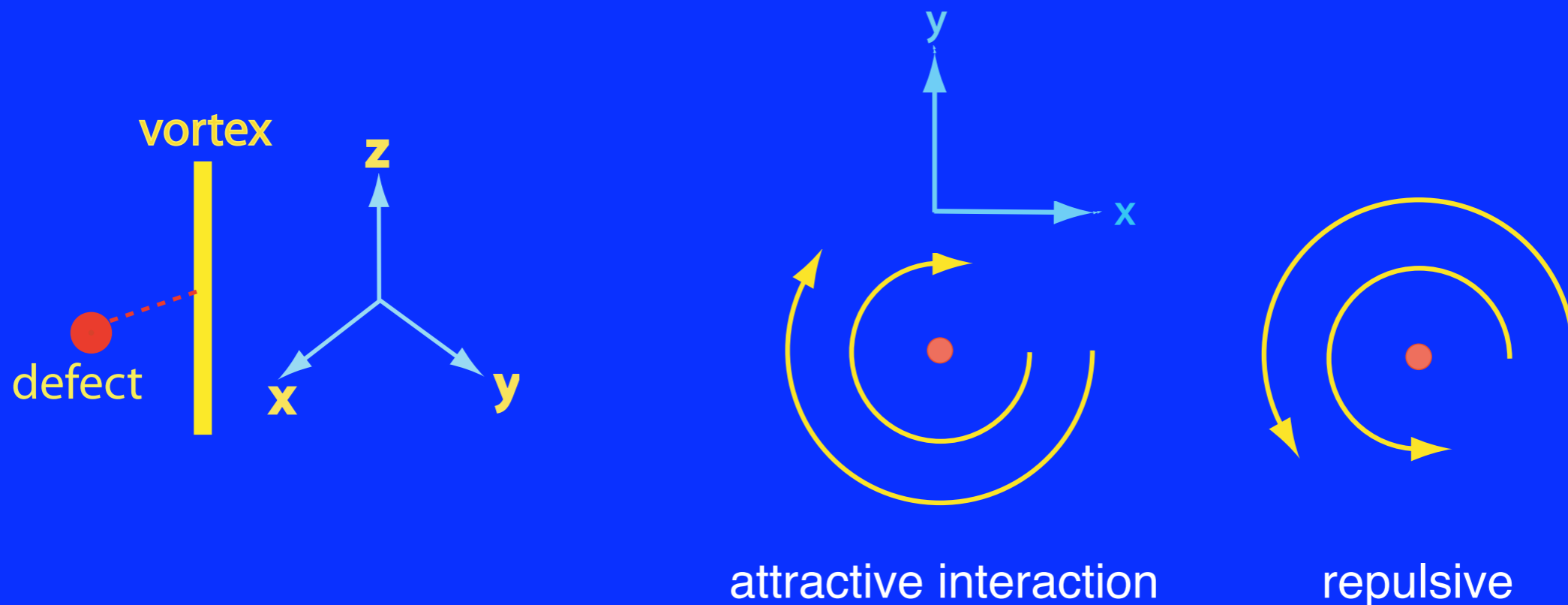
(A. Sedrakian, 95)

- Straight vortex segment.
- Single, point potential.



Motion of a segment

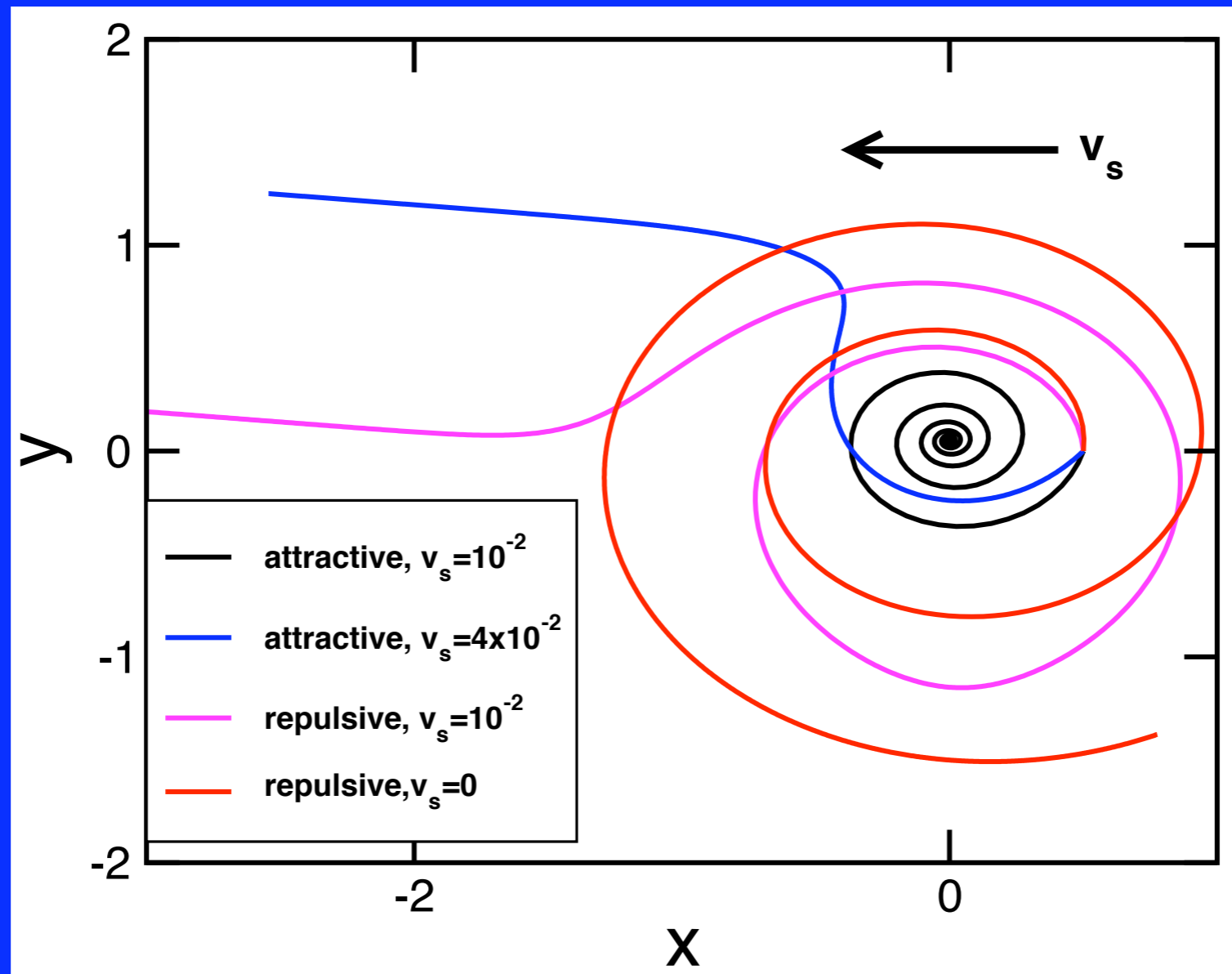
- Zero dissipation.
- SF flow = 0.



Vortex segment is trapped in an orbit

Motion of a segment

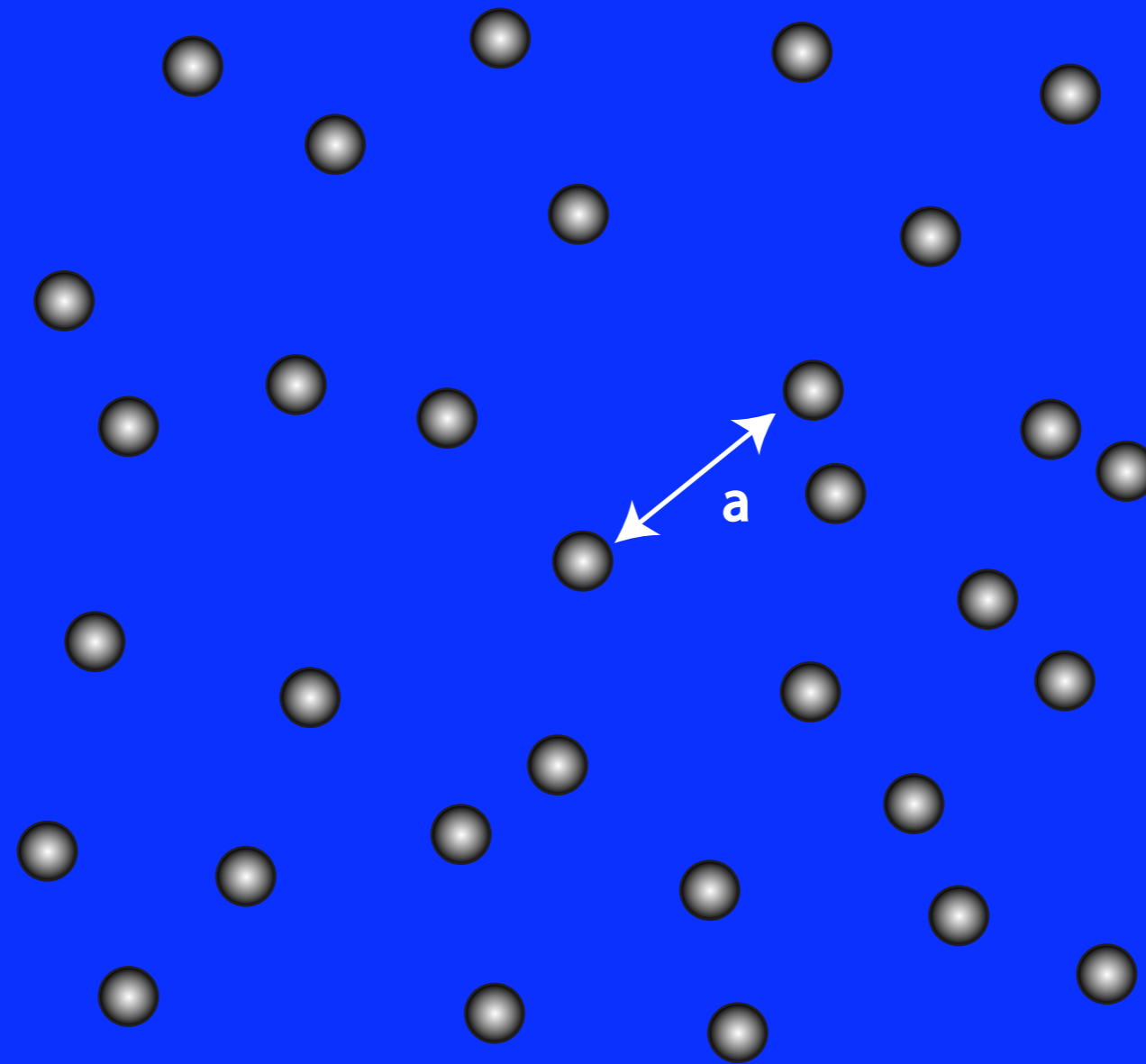
- with dissipation and
- ambient SF flow



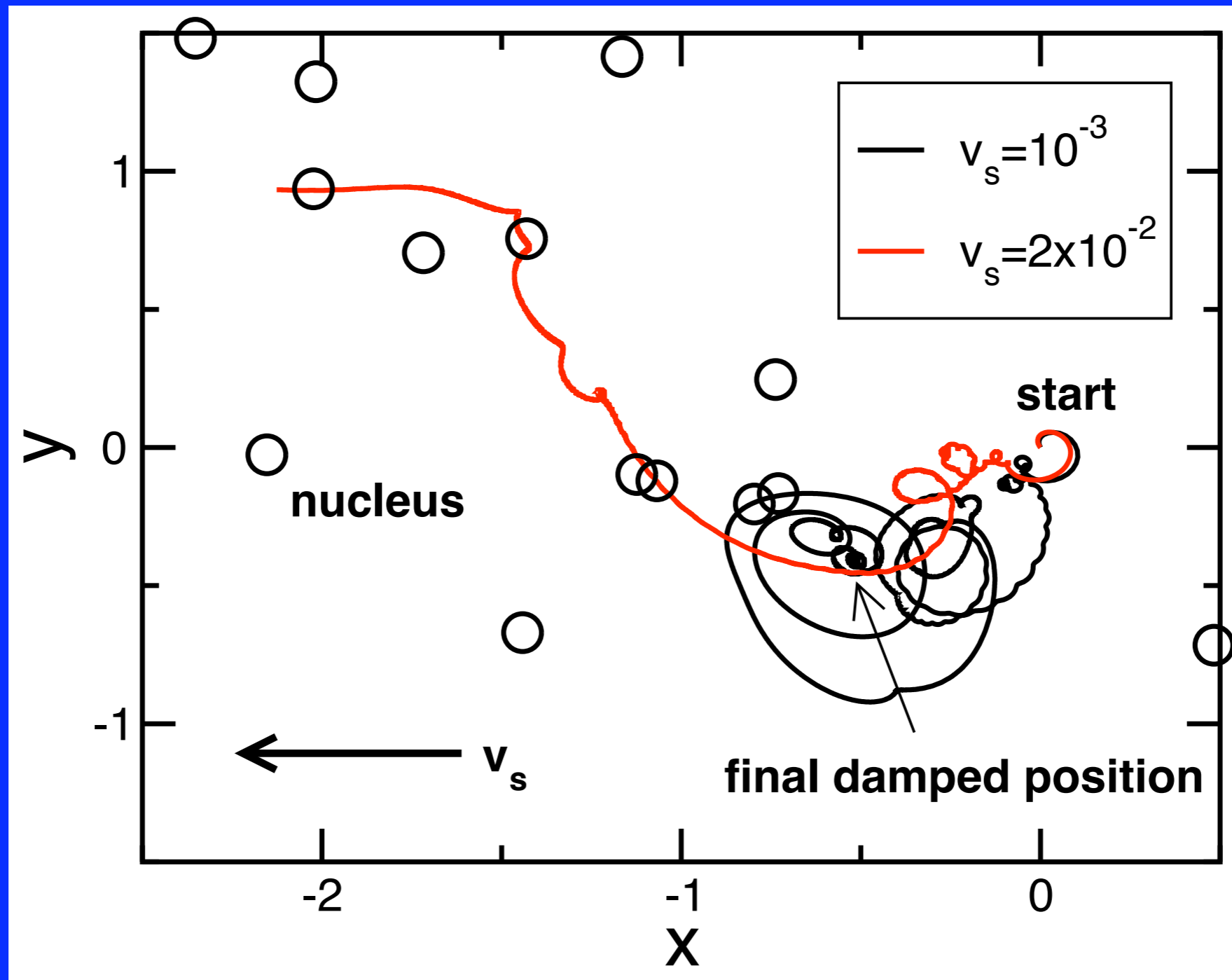
How does a vortex move
through the lattice?

The nuclear lattice is probably amorphous

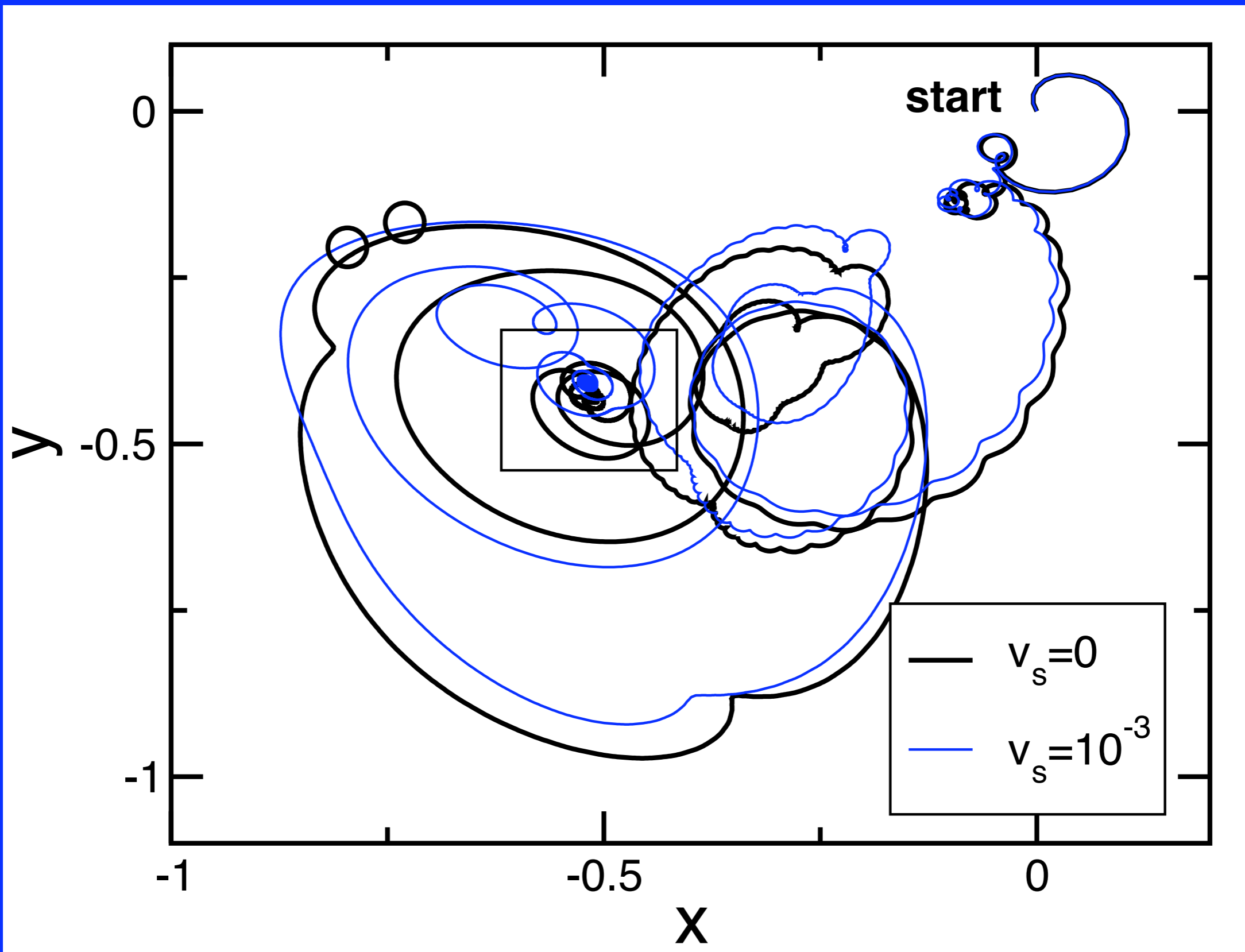
(Jones 01)



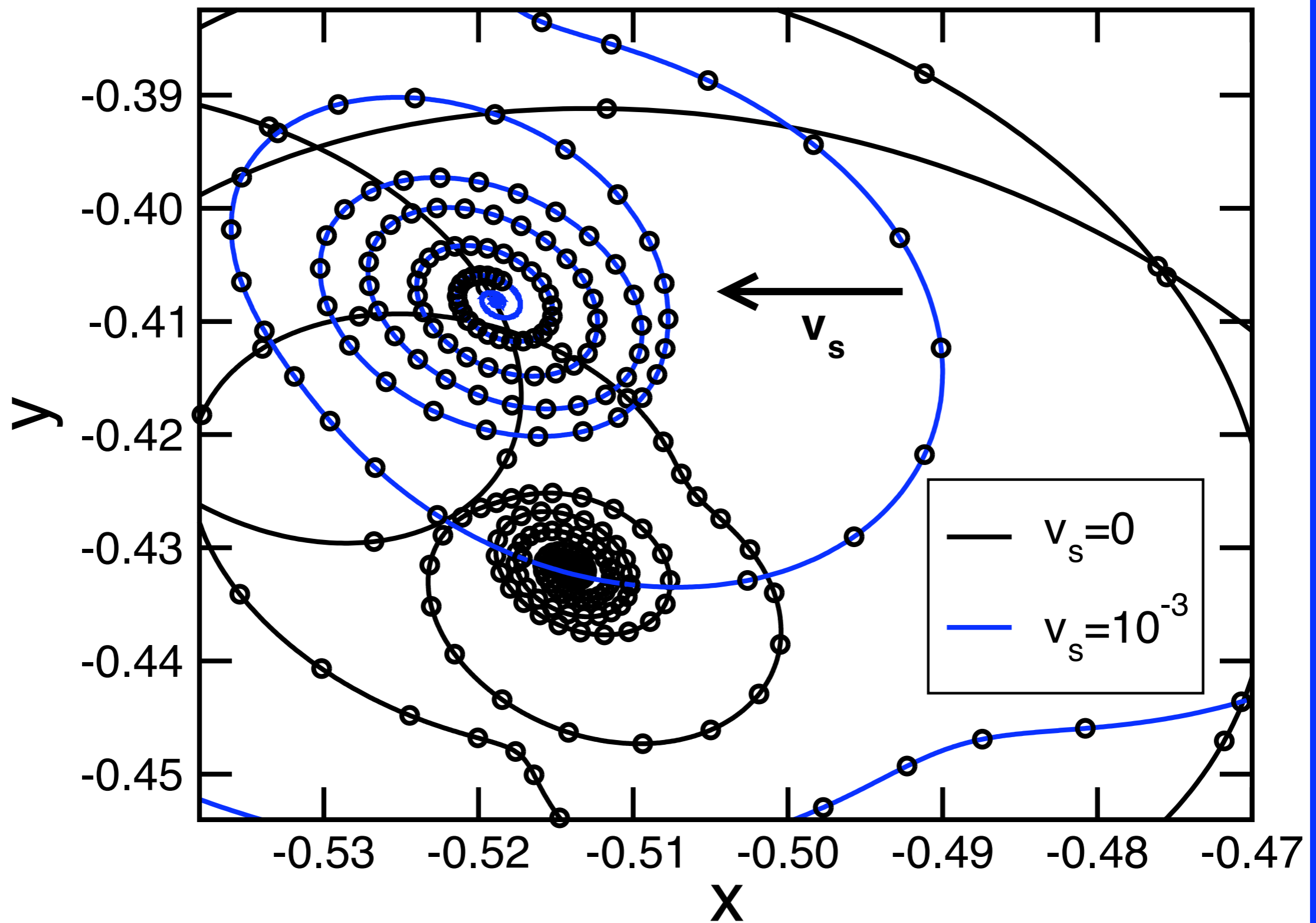
Motion in random lattice with dissipation: example segment



Pinning at low v_s

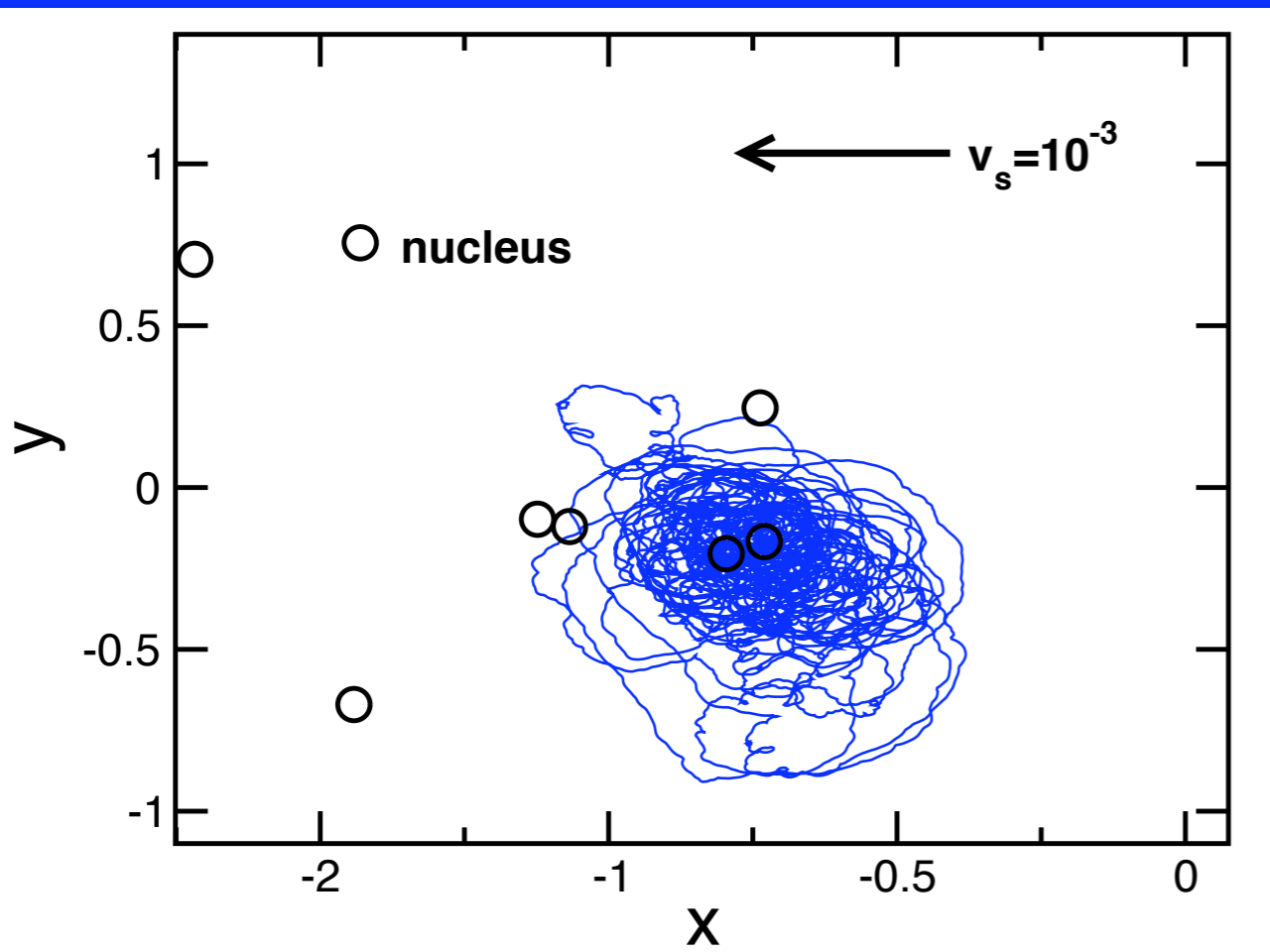


Pinning at low v_s (detail)

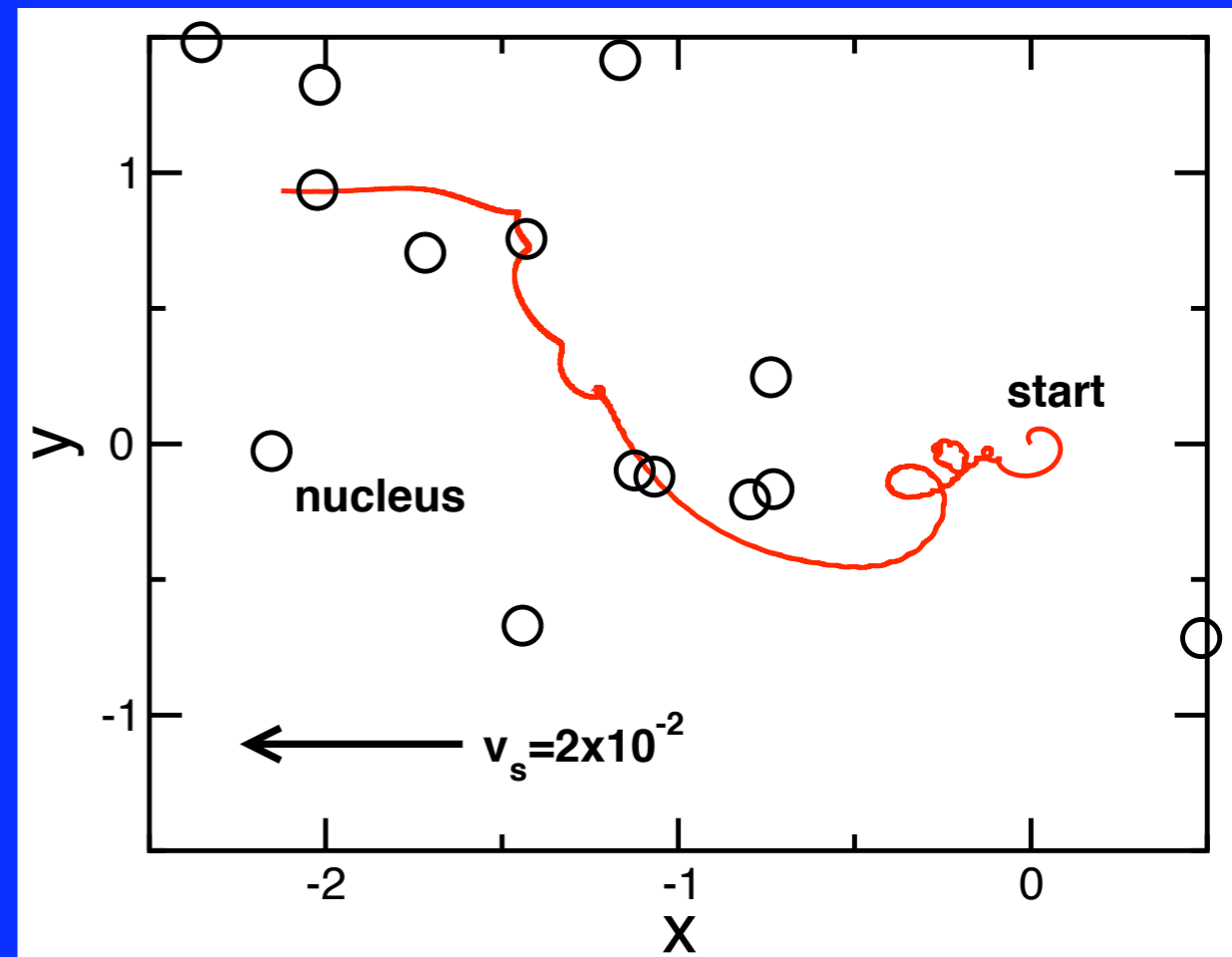


Motion with no dissipation

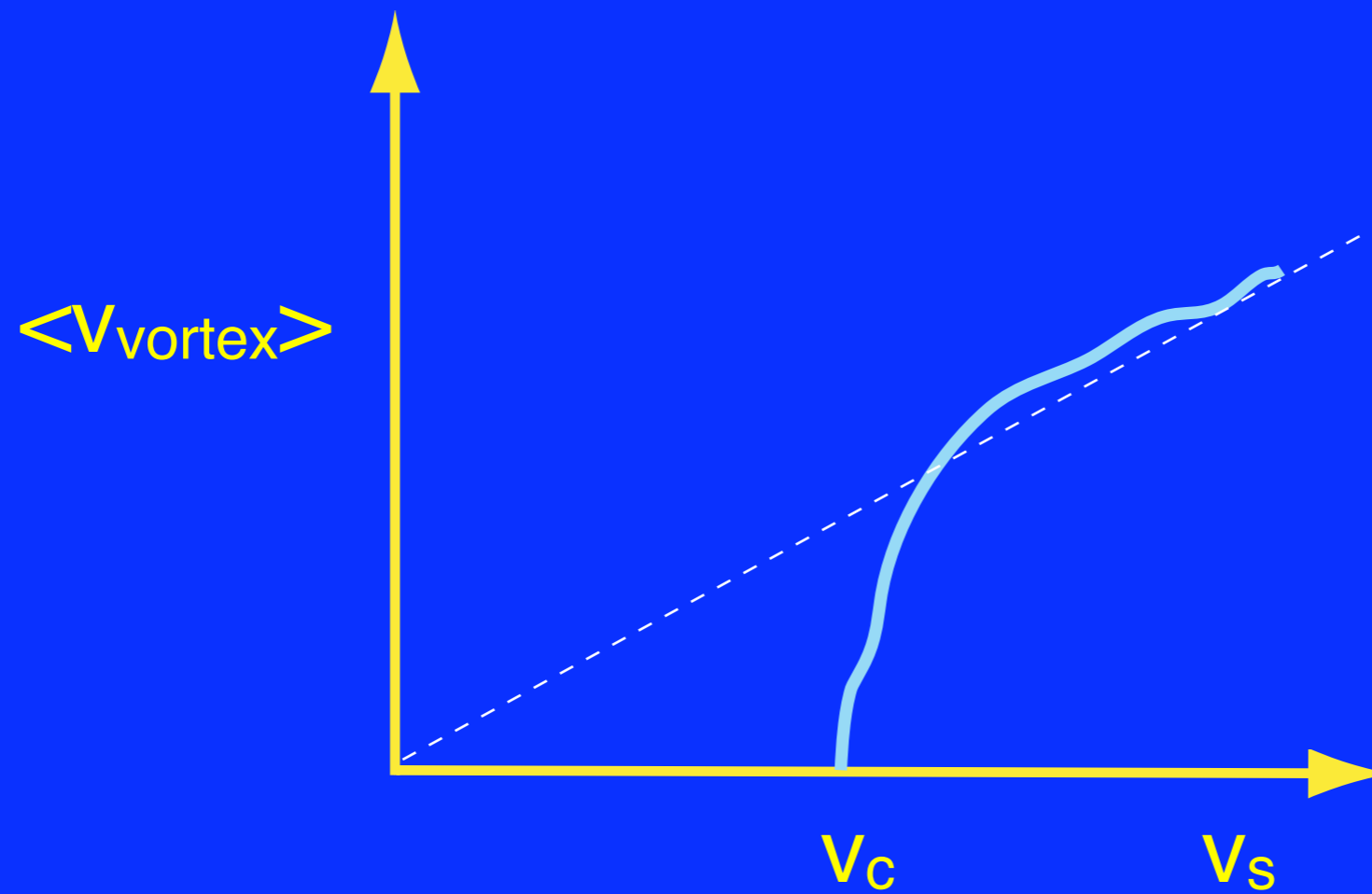
Below critical v_s



Above critical v_s



Existence of a transition



There is no translatory motion below a critical velocity,
even for zero drag.

The critical velocity

$$v_c \sim \frac{1}{\rho_s \kappa T_v^{1/2}} \left(\frac{E_{vn}}{a} \right)^{3/2} \sim \boxed{10^6 - 10^7 \text{ cm s}^{-1}}$$

E_{vn} = vortex-nucleus interaction energy

a = nuclear spacing

T_v = vortex tension

ρ_s = SF mass density

κ = vorticity quantum

Critical velocities this large (10^6 - 10^7 cm s⁻¹) can account for glitches

(Link & Cutler 02)

...but now there is another problem...

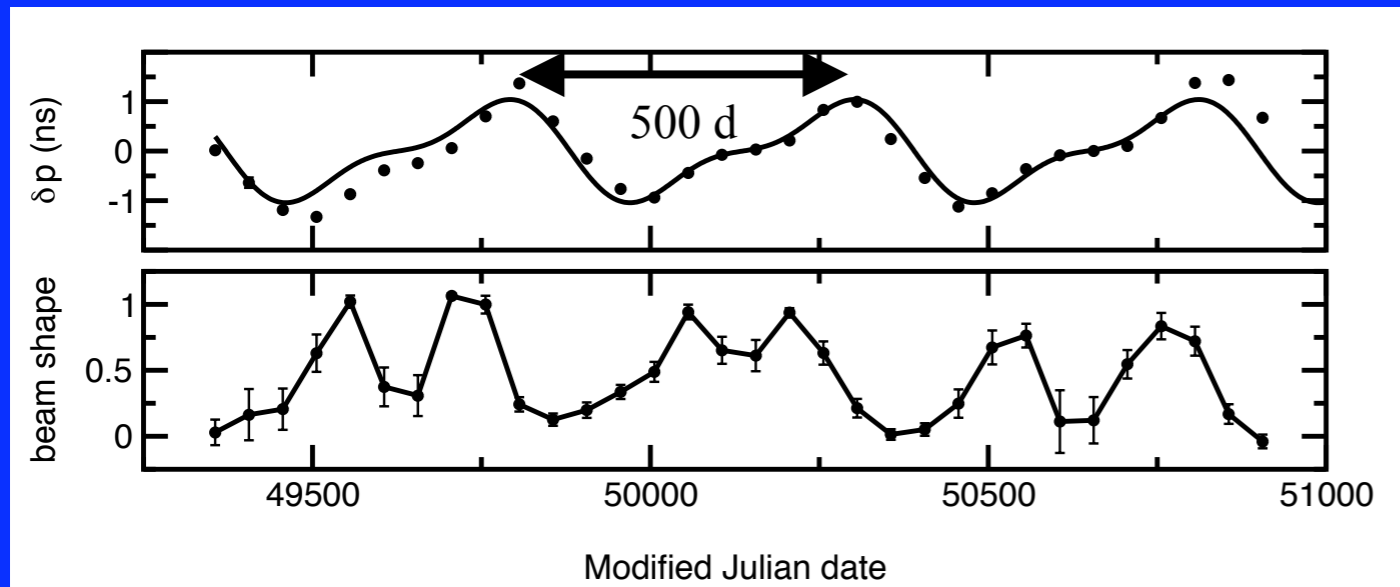
Some isolated neutron stars that appear to be precessing

- PSR 1828-11 (Stairs, Lyne, Shemar 00)
period of ~500 d.
- PSR B1642-03 (Shabanova, Lyne, Urama 01)
period of ~3 yr.
- RX J0720.4 (Haberl et al. 06)
period of ~7 yr.

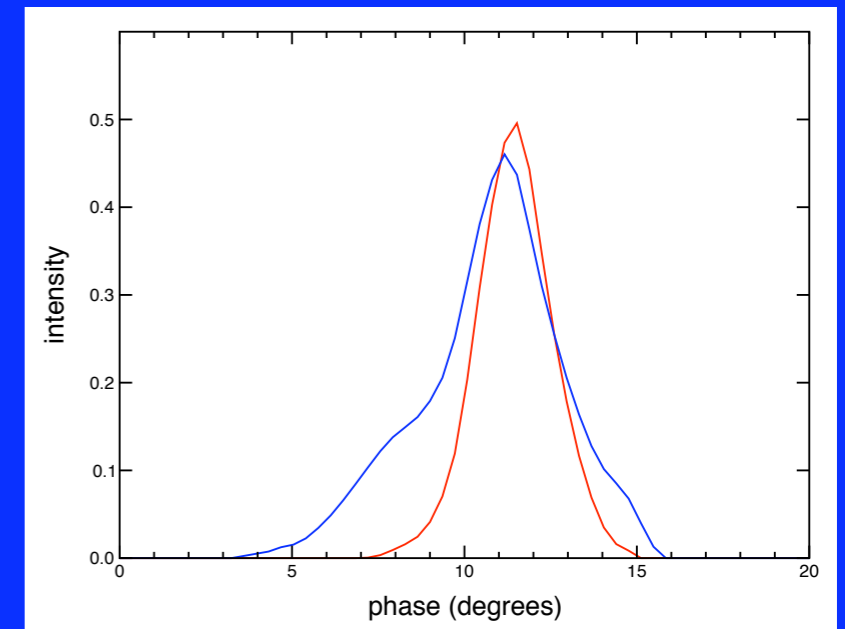
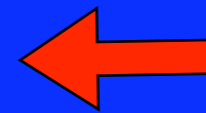
Evidence for precession:

B1828-11

(Stairs et al. 00)

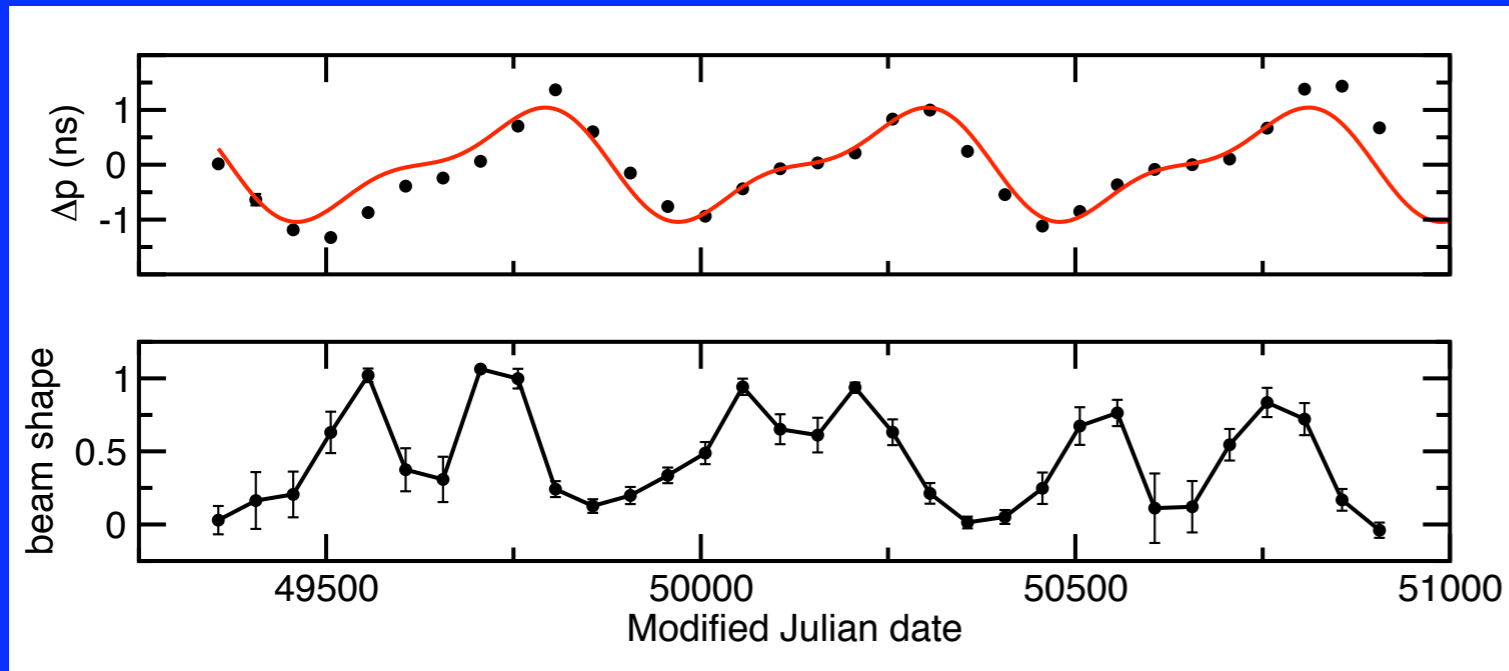


- Periodic but non-sinusoidal \Rightarrow harmonics.
- Strong periodicities at 511 and 256 d.
- Correlated changes in beam width.



Rigid-body modeling

(Link & Epstein 01; Akgun, Wasserman & Link 06)

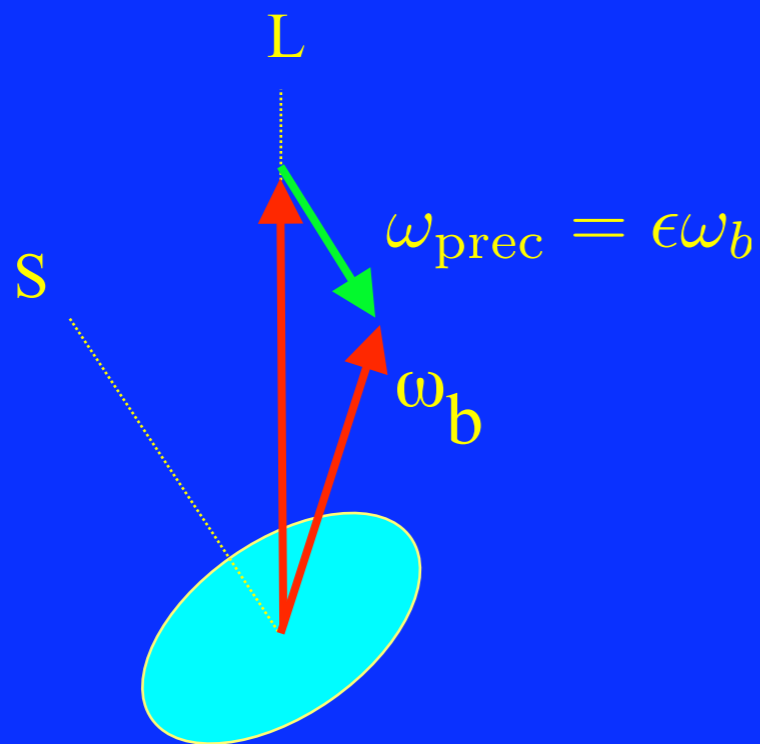


Wobble angle $\alpha = 3$ deg
Dipole angle $\chi =$ various
Precession period = 511 d
Harmonic at 256 d.

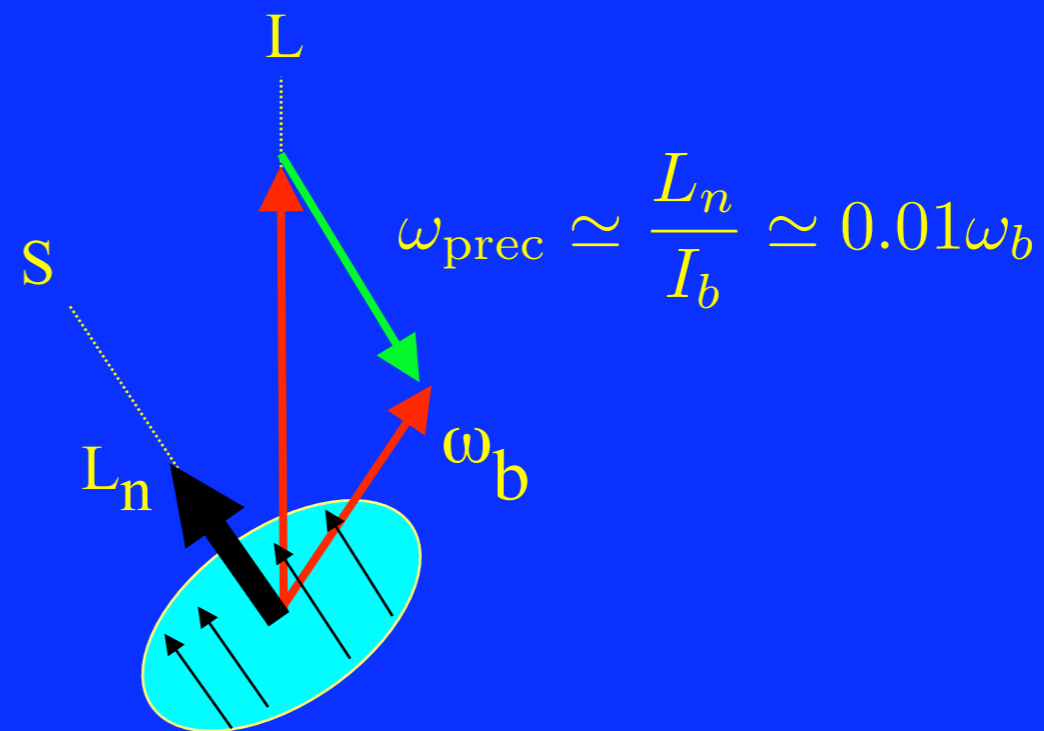
- Triaxiality (expected) gives a wide parameter space.
- Non-linearity in the torque or triaxiality produces the harmonic.

With immobilized vortices, precession is very fast

(Shaham 77; A. Sedrakian et al. 99; Link & Cutler 02)



Without pinning



With pinning
(gyroscopic effect)

PROBLEM: The precession frequency would be
 10^6 times faster than observed.

Conclusions

- Non-dissipative lattice force immobilizes inner-crust vortices.
- Inconsistent with long-period precession (in simplified, rigid-body hydrodynamics).
- Need a better understanding of the modes of rotating quantum liquid mixtures.



His inner core is spinning...

How does a neutron star do it?